

# Forced vibration of bending thick rectangular plates with different boundary conditions under concentrated load

CHEN Ying-jie<sup>1,2</sup>, YUE Xing-yu<sup>\*1,2</sup>

(1. College of Civil Engineering and Mechanics, Yanshan University, Qinhuangdao 066000, China;  
2. Hebei Province Civil Engineering Key Laboratory of Green Construction and Intelligent Operation  
and Maintenance, Qinhuangdao 066000, China)

**Abstract:** The boundary integral method is used to study the forced vibration of thick rectangular plates under three complex boundary conditions, namely, simply supported on four sides, fixed two opposite sides and simply supported on two opposite sides, and fixed on four sides. The solution process is clear, and the governing equations of forced vibration and the equations of the flexural surface are given. Through numerical calculation on the Matlab platform, the calculation results in the form of charts are obtained, and compared with the finite element simulation results. From the comparison of these results the accuracy of applying the boundary integral method to solve the forced vibration problem of thick rectangular plates and the correctness of the derived governing equations and the flexural surface equations can be seen. It has certain practical significance for various related problems in engineering practice, and provides a new way for solving such problems, which can be directly applied to engineering practice.

**Key words:** theory of thick plates; boundary integral method; the forced vibration; the boundary conditions; numerical calculation; the deflection surface equations

## 1 Introduction

Thick plate structure theory has been widely used in modern science and technology and production activities<sup>[1]</sup>. It not only has a large number of theoretical basis for the protection engineering and atomic energy engineering of thick plate components, but also needs to involve and apply the thick plate structure theory in modern engineering such as aviation, aerospace and ships<sup>[2]</sup>. At the same time, the forced vibration of thick plate is also inevitable, and the research on the vibration of thick plate is becoming more and more important<sup>[3]</sup>. However, since the thin plate does not consider the influence of shear deformation on bending, the classical thin plate theory cannot meet the accuracy requirements<sup>[4-7]</sup>. There-

fore, it is very necessary to propose a new method to solve the thick plate problem.

In the middle of the 20<sup>th</sup> century, Reissner<sup>[8-10]</sup> considered the influence of shear deformation and extrusion deformation, and deduced the static equation of thick plate problem. Mindlin<sup>[10,11]</sup> considered the influence of transverse shear deformation and rotational inertia, and deduced the dynamic equation of bending vibration of thick plates. Up to now, the thick plate theory proposed by Reissner and Mindlin is still widely used by modern scholars as the basic equation for studying thick plate problems.

Based on Mindlin plate theory, Wu et al.<sup>[12]</sup> solved the free vibration problem of plate with arbitrary shape by using differential cubature method. The typical examples of thick plate bending are calculated in detail based on Hamilton method<sup>[13]</sup>. Based on Mindlin plate theory, Li et al.<sup>[14]</sup> proposed a high-order eight-node hybrid stress quadrilateral element, and analyzed the bending and free

Received by: 2021-04-10; Revised by: 2021-09-20.

Corresponding authors: YUE Xing-yu\* (1998-), Female,  
Master Reading  
(E-mail: yxy9367@163.com).

vibration of simply supported, clamped square plates and circular plates with different thickness-span ratios.

Among the existing theories for solving thick plate components, the Reissner theory is the best, while this paper studies the forced vibration of thick rectangular plates by using the boundary integral method on the basis of the Reissner theory. This provides a new method for studying the forced vibration of thick rectangular plates.

## 2 Forced vibration of simply supported thick rectangular plate on four sides

### 2.1 Deflection surface equations

According to the governing equation<sup>[15]</sup>, the governing equation of forced vibration of the thick plate is

$$D\nabla^4 w = \left[ p\delta(x-x_0, y-y_0) - \frac{kh^2}{10} p\nabla^2 \delta(x-x_0, y-y_0) \right] + D\lambda^2 w - \frac{kh^2}{10} D\lambda^2 \nabla^2 w \quad (1)$$

where  $h$  is the thickness of the plate,  $k$  is the volume modulus, and  $D$  is the bending stiffness of the plate.

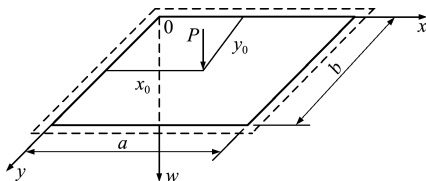


Fig. 1 Amplitude actual system of simply supported thick rectangular plate with four sides under concentrated load

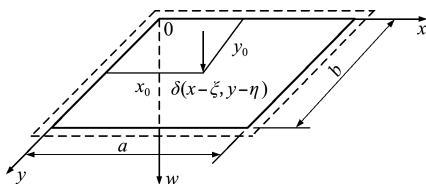


Fig. 2 Amplitude quasi-basic system of bending thick rectangular plate

The simply supported thick rectangular plate with four sides under a concentrated load shown in Fig.1 is taken as the amplitude actual system. The bending thick rectangular plate simply supported on four sides under a unit of concentrated load shown in Fig.2 is taken as the quasi-basic system.  $w_1$  represents the deflection of the basic system and  $w$  represents the deflection of the actual system.

The boundary integral method is applied between the amplitude actual system of simply supported four sides shown in Fig.1 and the quasi-basic system shown in Fig.2,

$$w(\xi, \eta) = \int_0^a \int_0^b \left[ p\delta(x-x_0, y-y_0) - (kh^2/10) P\nabla^2 \delta(x-x_0, y-y_0) + D\lambda^2 w(x, y) - \frac{kh^2}{10} D\lambda^2 \nabla^2 w(x, y) \right] w_1(x, y, \xi, \eta) dx dy \quad (2)$$

From Ref. [15], it can be seen that the expressions of  $w_1$  are

$$w_1(x, y; a-\xi, \eta) = -\frac{2}{Db} \sum_{n=1,2}^{\infty} \frac{1}{k_n^2 - \lambda_n^2} \times \left[ \frac{\sinh k_n(a-\xi)}{k_n \sinh k_n a} \sinh k_n x - \frac{\sinh \lambda_n(a-\xi)}{\lambda_n \sinh \lambda_n a} \sinh \lambda_n x \right] \cdot \sin \beta_n \eta \sinh \beta_n y \quad (0 \leq x \leq \xi) \quad (3)$$

$$w_1(a-x, y; \xi, \eta) = -\frac{2}{Db} \sum_{n=1,2}^{\infty} \frac{1}{k_n^2 - \lambda_n^2} \times \left[ \frac{\sinh k_n(a-x)}{k_n \sinh k_n a} \sinh k_n \xi - \frac{\sinh \lambda_n(a-x)}{\lambda_n \sinh \lambda_n a} \sinh \lambda_n \xi \right] \cdot \sin \beta_n \eta \sinh \beta_n y \quad (\xi \leq x \leq a) \quad (4)$$

The other two forms of deflection curve equation can be found in Ref.[15], here is no longer given.

The expressions of  $w_1$  are substituted into Eq.(2), and the deflection surface equations can be obtained by simplifying the calculation.

### 2.2 Stress function

Assuming that the stress function is

$$\varphi(\xi, \eta) = \sum_{n=0,1,2}^{\infty} [E_n \cosh \delta_n \xi + F_n \cosh \delta_n (a-\xi)] \times \cos \beta_n \eta + \sum_{m=0,1,2}^{\infty} [G_m \cosh \gamma_m \eta + H_m \cosh \gamma_m (b-\eta)] \cos \alpha_m \xi \quad (5)$$

$$\delta_n = \sqrt{\beta_n^2 + 10/h^2}, \quad \gamma_m = \sqrt{\alpha_m^2 + 10/h^2} \quad (6)$$

The stress function should satisfy  $\nabla^2 \varphi - \frac{10}{h^2} \varphi = 0$ , which will be proved by the following formula

$$\nabla^2 \varphi = \sum_{n=0,1}^{\infty} [E_n \cosh \delta_n \xi + F_n \cosh \delta_n (a-\xi)] \times (\delta_n^2 - \beta_n^2) \cos \beta_n \eta + \sum_{m=0,1}^{\infty} [G_m \cosh \gamma_m \eta + H_m \cosh \gamma_m (b-\eta)] (\gamma_m^2 - \alpha_m^2) \cos \alpha_m \xi =$$

$$(10/h^2)\varphi \quad (7)$$

It can be seen from Eq.(7) that the assumed stress function satisfies the condition.

The bending moments of the simply supported bending thick rectangular plate on four sides are 0, and the following equation can be obtained

$$E_n = F_n = G_m = H_m = 0 \quad (8)$$

At this time, only the deflection surface equations and the assumed stress function equation are substituted into the expression to solve the bending moment amplitude of the thick rectangular plate.

Substitute Eq.(8) into the stress function expression (5), and get

$$\varphi(\zeta, \eta) = E_0 \cosh \delta_0 \zeta + F_0 \cosh \delta_0 (a - \zeta) + G_0 \cosh \gamma_0 \eta + H_0 \cosh \gamma_0 (b - \eta) \quad (9)$$

Then, according to the four-side torsion angles of

the simply supported bending thick rectangular plate are zero,  $E_0 = F_0 = G_0 = H_0 = 0$ .

At this time, only the deflection surface equations and the assumed stress function equation are substituted into equations  $\omega_x$  and  $\omega_y$ . Therefore, the final expression of stress function is

$$\varphi(\xi, \eta) = 0 \quad (10)$$

### 2.3 Numerical calculation

Combined with practical engineering application, taking the basic parameters:  $x_0 = y_0 = 0.5 \text{ m}$  (i.e. concentrated load  $P$  acting on the center point of rectangular plate),  $a = b = 1 \text{ m}$ ,  $P = 100 \text{ N}$ ,  $E = 200 \text{ GPa}$ ,  $\nu = 0.3$ ,  $h/a = 0.1, 0.2, 0.3$ , the frequency range is  $\omega/\omega_0 = 0.1, 0.3, 0.5, 0.6$ . The deflection curves of simply supported rectangular plate on four sides are obtained by numerical calculation.

Tab. 1 Finite element method and the method in this paper are used to calculate the deflection values of a forced vibration plate with four sides simply supported ( $h/a = 0.1, x/a = 0.5$ )  $10^{-10} \text{ m}$

$y/b$	$0.1\omega_0$		$0.3\omega_0$		$0.5\omega_0$		$0.6\omega_0$	
	Ansys	This paper	Ansys	This paper	Ansys	This paper	Ansys	This paper
0	0	0	0	0	0	0	0	0
0.1	168.65	168.16	184.78	183.98	227.37	225.60	269.04	266.06
0.2	332.76	331.82	363.48	361.97	444.63	441.26	523.99	518.31
0.3	485.86	484.50	528.27	526.11	640.19	635.46	749.58	741.66
0.4	617.74	616.21	667.72	665.25	799.54	794.04	928.31	919.06
0.5	736.85	745.07	789.25	796.69	927.36	932.22	1062.26	1063.80

Tab. 2 Finite element method and the method in this paper are used to calculate the deflection values of a forced vibration plate with four sides simply supported ( $h/a = 0.2, x/a = 0.5$ )  $10^{-10} \text{ m}$

$y/b$	$0.1\omega_0$		$0.3\omega_0$		$0.5\omega_0$		$0.6\omega_0$	
	Ansys	This paper	Ansys	This paper	Ansys	This paper	Ansys	This paper
0	0	0	0	0	0	0	0	0
0.1	23.32	23.16	25.62	25.30	31.68	30.85	37.62	36.15
0.2	46.41	46.08	50.79	50.17	62.36	60.76	73.68	70.86
0.3	68.97	68.44	75.03	74.09	91.02	88.73	106.65	102.68
0.4	91.63	90.46	98.79	97.13	117.67	114.41	136.10	130.86
0.5	133.47	132.91	140.93	139.94	160.57	158.15	179.72	175.47

Tab. 3 Finite element method and the method in this paper are used to calculate the deflection values of a forced vibration plate with four sides simply supported ( $h/a = 0.3, x/a = 0.5$ )  $10^{-10} \text{ m}$

$y/b$	$0.1\omega_0$		$0.3\omega_0$		$0.5\omega_0$		$0.6\omega_0$	
	Ansys	This paper	Ansys	This paper	Ansys	This paper	Ansys	This paper
0	0	0	0	0	0	0	0	0
0.1	8.02	7.92	8.83	8.65	10.99	10.51	13.10	12.26
0.2	16.14	15.93	17.70	17.32	21.82	20.88	25.85	24.23
0.3	24.55	24.18	26.72	26.10	32.43	31.04	38.01	35.68
0.4	35.50	33.44	38.05	35.73	44.76	41.59	51.30	47.07
0.5	61.48	59.03	64.14	61.45	71.14	67.64	77.95	73.42

### 3 Numerical examples

#### 3.1 Deflection surface equations

Firstly, the bending thick rectangular plate with two opposite sides fixed and two opposite sides simply supported under a concentrated load is taken as numerical example ①, Fig.3 is taken as the amplitude actual system. The bending moment constraints of the two fixed sides are removed and replaced by the distributed bending moments  $M_{x0}$  and  $M_{xa}$ , so the equivalent diagram of the amplitude actual system shown in Fig.4, and assume

$$M_{x0} = \sum_{n=1,2}^{\infty} A_n \sin \beta_n y \quad (11)$$

$$M_{xa} = \sum_{n=1,2}^{\infty} B_n \sin \beta_n y \quad (12)$$

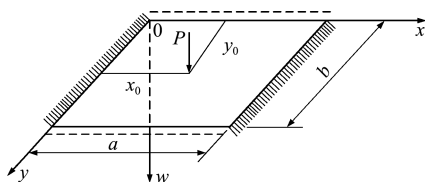


Fig. 3 Amplitude actual system of bending thick rectangular plate with two opposite sides fixed and two opposite sides simply supported

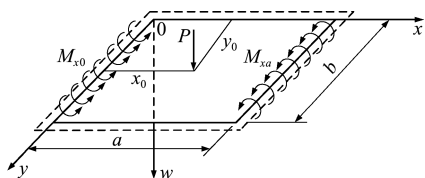


Fig. 4 Actual system equivalent diagram of the amplitude of bending thick rectangular plate with two opposite sides fixed and two opposite sides simply supported

In addition, a thick rectangular plate with four fixed sides under concentrated force is taken as numerical example ②. Assuming that the bending moments are

$$M_{x0} = \sum_{n=1,2}^{\infty} A_n \sin \beta_n y \quad (13)$$

$$M_{xa} = \sum_{n=1,2}^{\infty} B_n \sin \beta_n y \quad (14)$$

$$M_{y0} = \sum_{m=1,2}^{\infty} C_m \sin \alpha_m x \quad (15)$$

$$M_{yb} = \sum_{m=1,2}^{\infty} D_m \sin \alpha_m x \quad (16)$$

For thick rectangular plates with the above two boundary conditions, the boundary integral method is applied between the equivalent diagram of the corresponding amplitude actual system and the amplitude quasi-basic system. The expressions

of the deflection and boundary angles of the basic system given in Ref. [15] are substituted into the equation, and the deflection surface equations can be obtained after calculation and simplification.

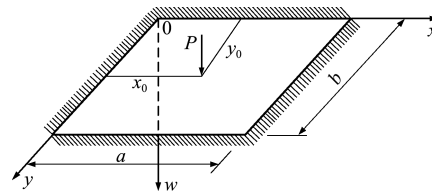


Fig. 5 Amplitude actual system of bending thick rectangular plate with fixed four sides

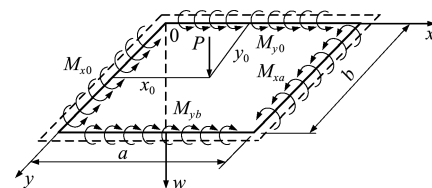


Fig. 6 Actual system equivalent diagram of the amplitude of bending thick rectangular plate with fixed four sides

#### 3.2 Stress function

For example ①, the bending thick rectangular plate with two opposite sides fixed and two opposite sides simply supported can be considered as a simply supported bending thick rectangular plate with four sides subjected to distributed bending moment respectively. From the foregoing, it can be seen that the  $E_0 = F_0 = G_0 = H_0 = 0$ ,  $G_m = H_m = 0$  of the four-sided simply supported thick rectangular plate, so the stress function here assumes that

$$\varphi(\zeta, \eta) = \sum_{n=1,2}^{\infty} [E_n \cosh \delta_n \xi + F_n \cosh \delta_n (a - \xi)] \cos \beta_n \eta \quad (17)$$

Since the internal bending moment caused by each deflection is balanced with the external bending moment, the stress function of this boundary condition can be obtained by calculation.

For example ②, the stress function is directly given here

$$\begin{aligned} \varphi(\zeta, \eta) = & \sum_{n=1,2}^{\infty} [B_n \cosh \delta_n \xi - A_n \cosh \delta_n (a - \xi)] \times \\ & [\beta_n / (\delta_n \sinh \delta_n a)] \cos \beta_n \eta - \\ & \sum_{m=1,2}^{\infty} [D_m \cosh \gamma_m \eta - C_m \cosh \gamma_m (b - \eta)] \times \\ & [\alpha_m / (\gamma_m \sinh \gamma_m b)] \cos \alpha_m \zeta \end{aligned} \quad (18)$$

#### 3.3 Boundary conditions

For example ①, the deflection surface equations and stress function should satisfy the following

boundary conditions

$$\omega_{\zeta\zeta 0} = \omega_{\zeta\zeta a} = 0 \quad (19)$$

The deflection surface equations and the stress function are substituted into the boundary conditions, and the boundary condition execution equations can be derived by calculation. Then the boundary execution equations are combined to obtain the values of unknown quantities  $A_n$  and  $B_n$ .

For numerical example ②, the following

boundary conditions should be satisfied to obtain the values of unknown quantities  $A_n, B_n, C_m$  and  $D_m$

$$\omega_{\zeta\zeta 0} = \omega_{\zeta\zeta a} = \omega_{\eta\eta 0} = \omega_{\eta\eta b} = 0 \quad (20)$$

### 3.4 Numerical calculation

For the above two examples, taking the size and mechanical parameters of the plate as the same as 2, 3, and the calculation results of example ① are listed in Tables 4 to 6, the numerical results of example ② are shown in Fig.7 to Fig.9.

Tab. 4 Finite element method and the method in this paper are used to calculate the deflection values of a forced vibration plate with two opposite sides simply supported, two opposite sides fixed ( $h/a=0.1, x/a=0.5$ )  $10^{-10} \text{ m}$

$y/b$	$0.1\omega_0$		$0.3\omega_0$		$0.5\omega_0$		$0.6\omega_0$	
	Ansys	This paper	Ansys	This paper	Ansys	This paper	Ansys	This paper
0	0	0	0	0	0	0	0	0
0.1	92.68	94.08	102.54	103.89	127.86	129.90	152.91	155.40
0.2	188.16	190.83	206.62	209.60	255.45	259.32	303.28	308.00
0.3	286.71	290.28	312.34	316.34	380.02	385.22	446.18	452.54
0.4	383.50	387.76	413.88	418.62	493.94	500.08	572.06	579.56
0.5	491.27	504.80	523.16	537.36	607.11	623.24	688.97	706.96

Tab. 5 Finite element method and the method in this paper are used to calculate the deflection values of a forced vibration plate with two opposite sides simply supported, two opposite sides fixed ( $h/a=0.2, x/a=0.5$ )  $10^{-10} \text{ m}$

$y/b$	$0.1\omega_0$		$0.3\omega_0$		$0.5\omega_0$		$0.6\omega_0$	
	Ansys	This paper	Ansys	This paper	Ansys	This paper	Ansys	This paper
0	0	0	0	0	0	0	0	0
0.1	14.68	14.86	16.24	16.36	20.37	20.28	24.43	24.01
0.2	29.97	30.29	32.96	33.16	40.87	40.66	48.63	47.80
0.3	46.33	46.69	50.50	50.70	61.51	61.12	72.26	71.02
0.4	65.00	64.88	69.97	69.65	83.04	82.02	95.78	93.74
0.5	105.80	106.01	110.98	111.05	124.61	124.13	137.87	136.50

Tab. 6 Finite element method and the method in this paper are used to calculate the deflection values of a forced vibration plate with two opposite sides simply supported, two opposite sides fixed ( $h/a=0.3, x/a=0.5$ )  $10^{-10} \text{ m}$

$y/b$	$0.1\omega_0$		$0.3\omega_0$		$0.5\omega_0$		$0.6\omega_0$	
	Ansys	This paper	Ansys	This paper	Ansys	This paper	Ansys	This paper
0	0	0	0	0	0	0	0	0
0.1	5.75	5.79	6.37	6.35	8.02	7.81	9.65	9.17
0.2	11.81	11.87	13.01	12.96	16.18	15.76	19.30	18.37
0.3	18.59	18.58	20.27	20.10	24.70	24.02	29.04	27.65
0.4	28.58	26.86	30.57	28.69	35.83	33.36	40.94	37.67
0.5	54.29	52.11	56.39	54.05	61.89	59.00	67.23	63.57

For the above two examples, the derivation process of the deflection surface equations, the stress function and the boundary execution equa-

tions is similar to that of the thick rectangular plate simply supported on four sides. Therefore, due to the limitation of space, it is no longer speci-

fied in this chapter.

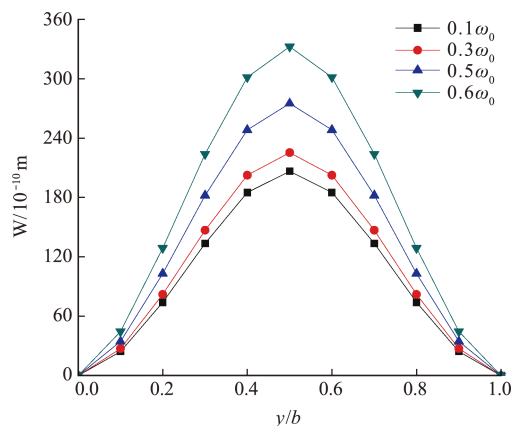


Fig. 7 Method in this paper is used to calculate the deflection values of a forced vibration plate with four sides fixed ( $h/a=0.1, x/a=0.3$ )

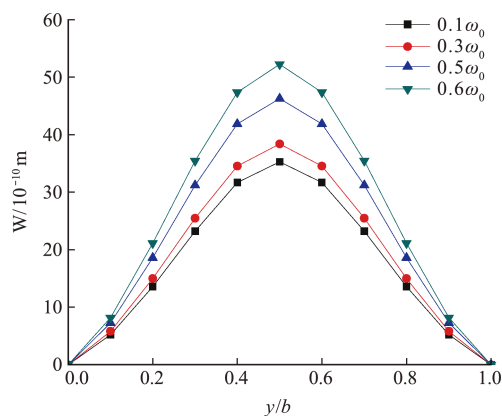


Fig. 8 Method in this paper is used to calculate the deflection values of a forced vibration plate with four sides fixed ( $h/a=0.2, x/a=0.3$ )

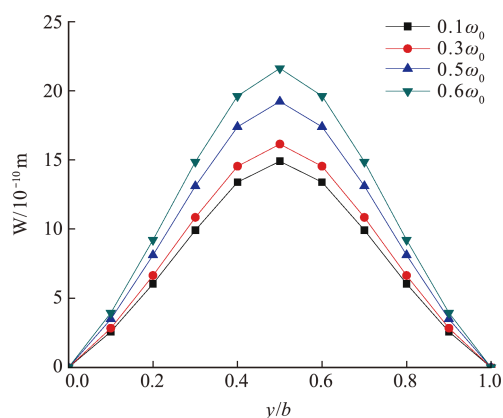


Fig. 9 Method in this paper is used to calculate the deflection values of a forced vibration plate with four sides fixed ( $h/a=0.3, x/a=0.3$ )

## 4 Result analysis

For simply supported thick rectangular plates with four sides, the comparison of the data from Tables 1 to 3 shows that when the frequency is

small and the plate thickness is small, the difference between the calculation results of this paper method and the finite element method is small. With the increase of thickness and vibration frequency, the relative error increases, but they are basically controlled within 5%, which has a certain impact on the error of the force acting point and its nearby nodes. However, for thick rectangular plates with two opposite sides fixed and two opposite sides simply supported, it can be seen from the comparison of data in Tab.4 to Tab.6 that the relative errors are basically controlled within 5%, and there is a large error at the action point of force and its adjacent nodes.

In addition, it can be seen from Fig.7 to Fig.9 that when the action frequency is small, the amplitude of the thick plate changes little. With the increase of frequency, the vibration amplitude of the thick rectangular plate increases sharply, which is also in line with the basic law of vibration.

## 5 Conclusion

Solving the thick plate problem is ultimately required to solve a higher order differential equation. In this paper, the boundary integral method is successfully applied to avoid this problem, so that it only needs a simple integration to obtain the deflection surface equations of the forced vibration of thick rectangular plates. Through numerical calculation on Matlab platform and simulation by finite element software, it can be seen that the boundary integral method is feasible to solve the forced vibration problem of thick rectangular plates, and the calculation results in this paper are correct. Therefore, it can be directly applied to engineering practice to solve some related engineering practical problems.

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## 集中载荷作用下不同边界条件的 弯曲厚矩形板的受迫振动

陈英杰<sup>1,2</sup>, 岳星宇<sup>\*1,2</sup>

(1. 燕山大学 建筑工程与力学学院, 秦皇岛 066000;

2. 河北省土木工程绿色建造与智能运维重点实验室, 秦皇岛 066000)

**摘要:**运用边界积分法研究了四边简支、两对边固定另两对边简支、四边固定三种复杂边界条件下厚矩形板的受迫振动问题,求解过程清晰,从而给出了受迫振动控制方程和挠曲面方程。通过在 Matlab 平台上进行数值计算,得出了图表形式的计算结果,并与有限元模拟值进行对照。研究表明,边界积分法用于求解厚矩形板的受迫振动问题的准确性,本文推导的控制方程和挠曲面方程的正确性,进而对工程实际中的各种相关问题具有一定的现实意义,也为求解此类问题提供了一种新途径,可以直接运用到工程实际中。

**关键词:**厚板理论;边界积分法;受迫振动;边界条件;数值计算;挠曲面方程

**中图分类号:**O32;O302

**文献标志码:**A

**文章编号:**1007-4708(2022)06-0845-07

收稿日期:2021-04-10;修改稿收到日期:2021-09-20.

作者简介:岳星宇\*(1998-),女,硕士生(E-mail:xyx9367@163.com).

引用本文/Cite this paper:

陈英杰,岳星宇. 集中载荷作用下不同边界条件的弯曲厚矩形板的受迫振动[J]. 计算力学学报, 2022, **39**(6):845-851.

CHEN Ying-jie, YUE Xing-yu. Forced vibration of bending thick rectangular plates with different boundary conditions under concentrated load[J]. *Chinese Journal of Computational Mechanics*, 2022, **39**(6):845-851.